Group Meeting #2

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NTHU, Interdisciplinary Program of Sciences

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Outline

The SIS Model

- Introduction to the SIS model.
- Non-equilibrium phase transition.
- Provide the TN approach
 - From Ising model to the SIS model.
 - The tensor product formulation.
 - The large deviation principle.



Introduction to the SIS Model



Introduction to the SIS Model

• The SIS model can be written as a differential equations system form.

$$\begin{cases} \frac{dS}{dt} = -\beta SI + \gamma I, \\ \frac{dI}{dt} = \beta SI - \gamma I. \end{cases}$$

- There are several kinds of extensions to the SIS model, for example, the SIR model.
- The SIS model is the simplest epidemic model that can describe the phase transitions.



Non-Equilibrium Phase Transition

The SIS model has the differential equation form

$$\begin{cases} \frac{dS}{dt} = -\beta SI + \gamma I, \\ \frac{dI}{dt} = \beta SI - \gamma I. \end{cases}$$

- 2 The model is characterized by two parameters β and γ .
- The non-equilibrium phase transition in SIS models is described by the relation

 $\begin{cases} \text{Absorbing Phase: } \frac{\beta}{\gamma} \leq 1, \\ \text{Endemic Phase: } \frac{\beta}{\gamma} > 1. \end{cases}$

Simulation:

https://kikiyenhaoyang.github.io/kikiyen/Web/TN_ftm

From Ising model to the SIS mode



Figure: One-dimensional Ising model.



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From Ising model to the SIS mode

● Each lattice in the model having the quantum state |s_i⟩ lives in the Hilbert space and therefore can be written as the linear combination of |0⟩ and |1⟩.

2 Composite Systems Postulate

The state space of a composite quantum system is the tensor product of the state spaces of its components. That is, the complete Hilbert space of a composite system is

$$\mathcal{H} = \bigotimes_{i=1}^{N} \mathcal{H}_i.$$



From Ising model to the SIS mode

The basis for the complete Hilbert space is

$$\{\bigotimes_{j=1}^N |i_j\rangle\}_{i_j\in\{0,1\}}.$$

Interpretation of the complete light model can be written as

$$|\Psi\rangle = \Psi(s_1, \cdots, s_N) \bigotimes_{i=1}^N |s_i\rangle.$$

From Ising Model to the SIS Model

- **(**) The Hibert space has dimension $2^N = \mathcal{O}(2^N)$, which grows rapidly.
- **2** With the TN, we can reduce it to $\mathcal{O}(N)$, which is much more efficient.
- The SIS model is also a kind of binary model. Thus we can write it in tensor product formulation as well.



Measurement Postulate

Upon measurement, the outcome of an observable is one of its eigenvalues, and the probability of obtaining a particular outcome is given by the square of the absolute value of the projection of the state vector onto the corresponding eigenvector. That is,

$$\sum_{i} C_{i} \ket{\Psi_{i}} \xrightarrow{\text{Measurement}} \ket{\Psi_{i}}, \text{ with probability } |C_{i}|^{2}$$



• In the lattice form SIS model. A lattice is in state $|0\rangle = |S\rangle$ or $|1\rangle = |I\rangle$ with probability $P_0(t)$ and $P_1(t)$ respectively. That is,

$$\left|P(t)
ight
angle=P_{0}(t)\left|0
ight
angle+P_{1}(t)\left|1
ight
angle,$$

which is analog to

$$\ket{\Psi} = \Psi_0 \ket{0} + \Psi_1 \ket{1}.$$

The state space of the SIS model is

$$\bigotimes_{i=1}^{N} \mathcal{H}_{i}, ext{ where } \mathcal{H} = \mathsf{Span}\{\ket{0}, \ket{1}\}.$$



• Since the SIS model is a kind of Markov chain, we can consider the infinitesimal Markov generator \hat{W} so that

$$\partial_t \left| P(t) \right\rangle = \hat{W} \left| P(t) \right\rangle,$$

which is Schrödinger-equation-like.

Itime Evolution Postulate

The time evolution of a quantum system is governed by the Schrödinger equation

$$-i\partial_t |\Psi\rangle = \hat{H} |\Psi\rangle,$$

which describes how the state vector changes over time.



 $\textbf{0} \ \ \text{The infinitesimal Markov generator} \ \ \hat{\mathcal{W}} \ \ \text{for the SIS model is}$

$$\hat{W} = \beta \sum_{i=1}^{N-1} (\hat{n}_i \omega_{i+1}^{\circ})^{\bullet \to 1} + \hat{\omega}_i^{\circ \to 1} \hat{n}_{i+1}) + \gamma \sum_{i=1}^{N} \hat{\omega}_i^{\circ \to 0} + \hat{W}_{driv}(\alpha),$$

where

$$\hat{n}_{i} = |1\rangle_{i} \langle 1|, \ \hat{\omega_{i}}^{0 \to 1} = |1\rangle_{i} \langle 0| - |0\rangle_{i} \langle 0|, \ \hat{\omega_{i}}^{0 \to 1} = |0\rangle_{i} \langle 1| - |1\rangle_{i} \langle 1|.$$

2 The generator is *analog* to the quantum Hamiltonian.



Quantum State Postulate

A quantum system is fully described by its state vector, usually denoted by $|\psi\rangle$ in Dirac notation. This state vector resides in a complex vector space known as a C–Hilbert space.



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Definition

Vector Space

Let F be a field. A F-vector space is set V equipped with two operators

$$+: V \times V \to V,$$

 $\cdot: F \times V \to V$

satisfies the following conditions

- $a \cdot (b \cdot v) = (ab) \cdot v$ for all $a, b \in F$ and $v \in V$,
- $(a+b) \cdot v = a \cdot v + b \cdot v$ for all $a, b \in F$ and $v \in V$,

•
$$a(u + v) = au + av$$
 for all $a \in F$ and $u, v \in V$,

• $id_F \in F$ such that $id_F \cdot v = v$.

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- Though we analog the generator to the quantum Hamiltonian, there are many differences between the generator and the quantum Hamiltonian.
- **②** The generator is a **REAL** matrix and the basis lives in a \mathbb{R} -Hilbert space.
- Many good properties are broken in the real Hilbert space.
- We can not treat it as a quantum system.

